

Solving the Bioheat Equation for Transcutaneous Recharging of a Medical Device Using Electric Fields

Susannah Engdahl

Senior Seminar Presentation

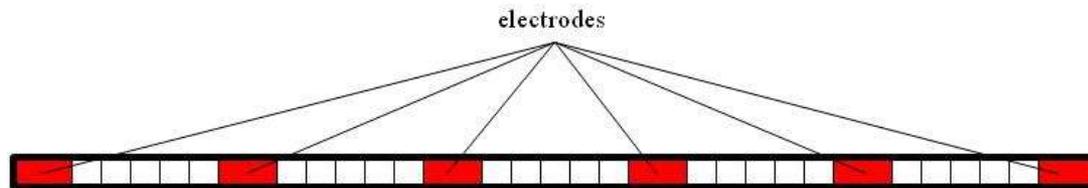
April 24, 2013

Electrophysiology

- The study of the body's electric activity
 - Can be small-scale (individual cells) or large-scale (entire organs)
- Electrophysiology often plays an important role in medical diagnostic procedures
 - Ex: ECG, EEG, EMG
- Signals often recorded by placing a series of electrodes on the surface of a patient's skin
 - Not always a practical approach—long-term collection of data may be required

A Possible Solution

- A subcutaneous (under the skin) recording device could remain in place semi-permanently
 - Device may be implanted almost anywhere in a minor surgical procedure



Schematic of proposed device

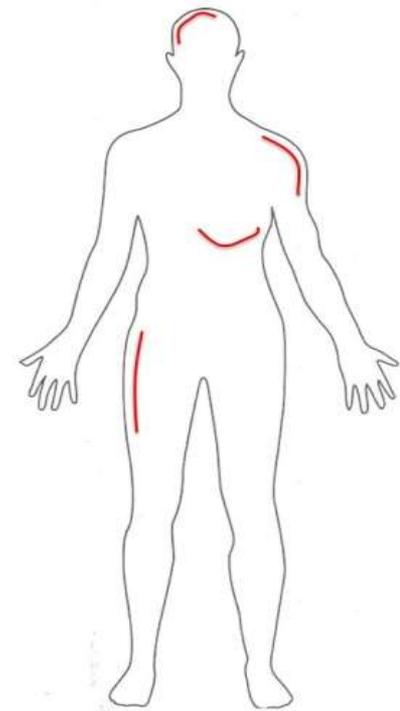
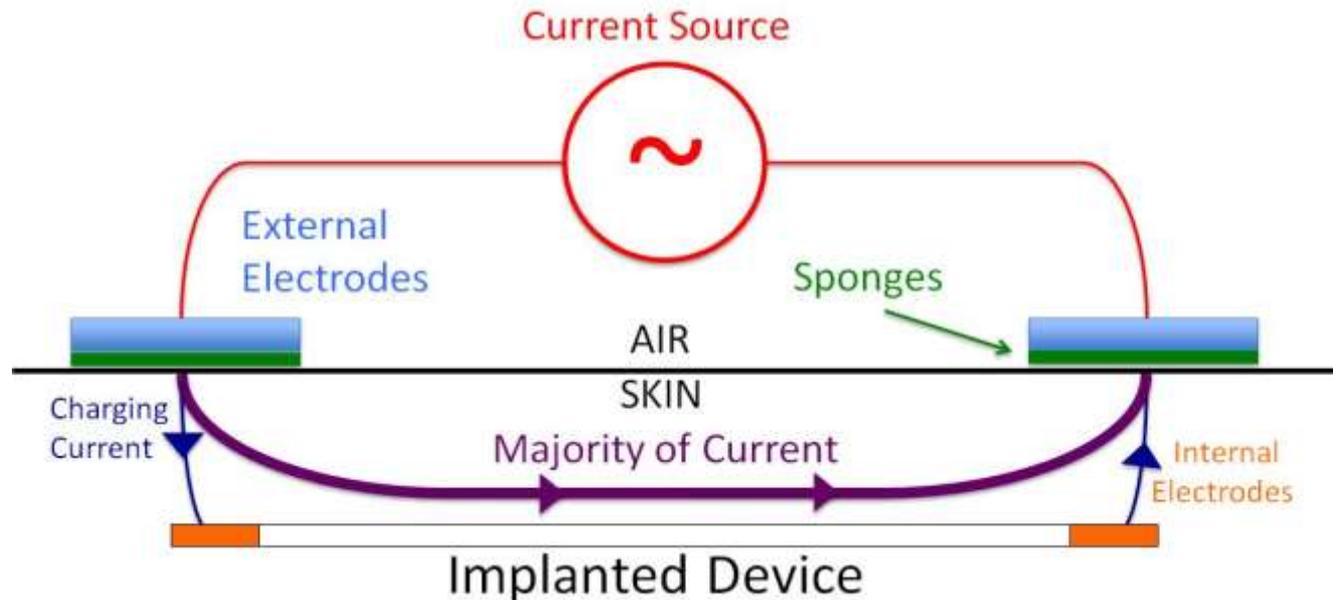


Figure created by Zachary Abzug

Transcutaneous Recharging

- Most implanted devices recharged via magnetic fields—not feasible for this device
- Instead, induce high frequency electric field using external source and sink electrodes



Transcutaneous recharging using electric fields

Figure created by Zachary Abzug

Upholding Medical Standards

- Problem with recharging via electric fields: current passing through tissue can cause thermal damage (Joule heating)
 - A temperature increase ≤ 2 °C is within medical standards
- Previous work: perform finite element analysis to investigate expected temperature increase

Project Objective

- Derive a closed-form solution for the anticipated temperature increase
 - Primary motivation: improved understanding of physical parameters on temperature increase

The “Extended” Bioheat Equation

$$\rho C \frac{\partial T}{\partial t} = k \nabla^2 T + \rho_b C_b \omega_b (T_b - T) + Q_{met} + J \cdot E$$



- The heat equation: describes variation of temperature in a region as a function of time
- Pennes’ extension to the heat equation: accounts for heat transfer due to perfusion (blood flow) and metabolic heat production
- Source term: accounts for heating due to power dissipation in tissue

Power Dissipation

- Calculate the work done by electromagnetic forces on a charge Q moving some infinitesimal distance $d\mathbf{l}$:

$$dW = \mathbf{F} \cdot d\mathbf{l} = Q(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \mathbf{v} dt$$

$$dW = Q\mathbf{E} \cdot \mathbf{v} dt$$

- Make the substitutions $Q = \int_V \rho d\tau$ and $\mathbf{J} = \rho\mathbf{v}$
- Rate at which work is done on all charges in a volume: $\frac{dW}{dt} = \int_V (\mathbf{J} \cdot \mathbf{E}) d\tau$
- Therefore, $\mathbf{J} \cdot \mathbf{E}$ is the work done on a charge per unit volume per unit time

Simplifying the Bioheat Equation

- Must make several simplifications

$$\rho C \frac{\partial T}{\partial t} = k \nabla^2 T + \rho_b C_b \omega_b (T_b - T) + Q_{met} + J \cdot E$$

- Steady-state solution ($\frac{\partial T}{\partial t} = 0$)
- Ignore perfusion ($\omega_b = 0$)
- Ignore metabolic heat production ($Q_{met} = 0$)
- Final equation to solve:

$$\nabla^2 T = -\frac{1}{k} (J \cdot E)$$

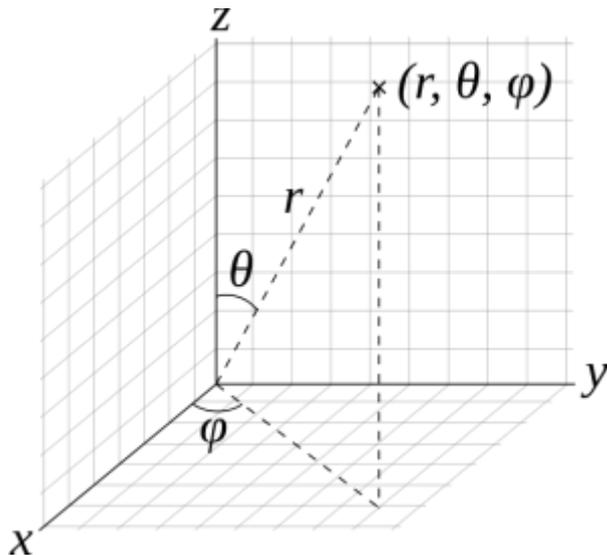
A Previous Solution

- Our method of solving the bioheat equation is similar to a solution given by Elwassif et al. [1]
 - Begin by relating the source term to the gradient of the electric potential

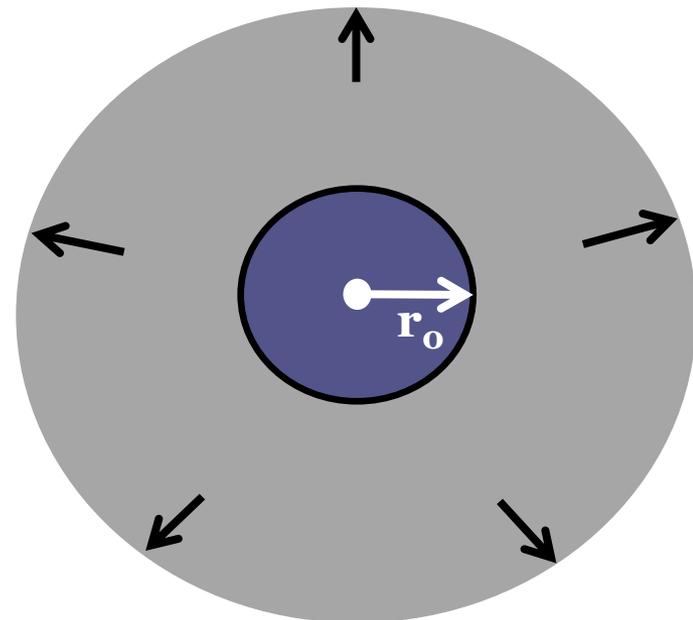
$$J \cdot E = \sigma E \cdot E = -\sigma \nabla V \cdot -\nabla V = \sigma |\nabla V| |\nabla V| \cos\theta = \sigma |\nabla V|^2$$

Geometric Considerations

- Treat electrode as a current-producing sphere in an infinite homogeneous and isotropic resistive material
- Second electrode is at infinity ($V=0$)



Spherical coordinate system
(image from Wikipedia)



Spherical electrode within infinite medium

A Solution in Spherical Coordinates

- Write bioheat equation in spherical coordinates
 - Ignore θ and φ dependence due to geometry

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial T}{\partial r} \right) = -\frac{\sigma}{k} |\nabla V|^2$$

- Must find an expression for the electric potential by solving Laplace's equation

Solving Laplace's Equation

- Laplace's equation is $\nabla^2 V = 0$
 - Write in spherical coordinates: $\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) = 0$
 - Multiply through by r^2 : $\frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) = 0$
 - Integrate twice with respect to r : $V(r) = -\frac{A}{r} + B$
 - Since zero potential at infinity, $B=0$: $V(r) = -\frac{A}{r}$

Solving Laplace's Equation

- To determine A, consider a point source of current in an infinite, homogeneous, isotropic medium.

- The current density is: $J = \frac{I}{4\pi r^2} \hat{r}$

- Since $J = \sigma E = -\sigma \nabla V$, the potential is:

$$\nabla V = -\frac{I}{4\pi\sigma r^2} \hat{r} \quad \longrightarrow \quad \frac{dV}{dr} = -\frac{I}{4\pi\sigma r^2} \quad \longrightarrow \quad V(r) = \frac{I}{4\pi\sigma r}$$

- Compare to $V(r) = -\frac{A}{r} \quad \longrightarrow \quad A = -\frac{I}{4\pi\sigma}$

Solving the Bioheat Equation

- Plug solution for potential into bioheat equation:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial T}{\partial r} \right) = -\frac{\sigma}{k} \left| \nabla \left(-\frac{A}{r} \right) \right|^2$$

product rule for derivatives

take gradient with respect to r

$$\frac{2}{r} \frac{\partial T}{\partial r} + \frac{\partial^2 T}{\partial r^2} = -\left(\frac{l}{4\pi} \right)^2 \frac{1}{k\sigma r^4}$$

$$2r^3 \frac{\partial T}{\partial r} + r^4 \frac{\partial^2 T}{\partial r^2} = -\left(\frac{l}{4\pi} \right)^2 \frac{1}{k\sigma}$$

- Make substitutions for clarity:

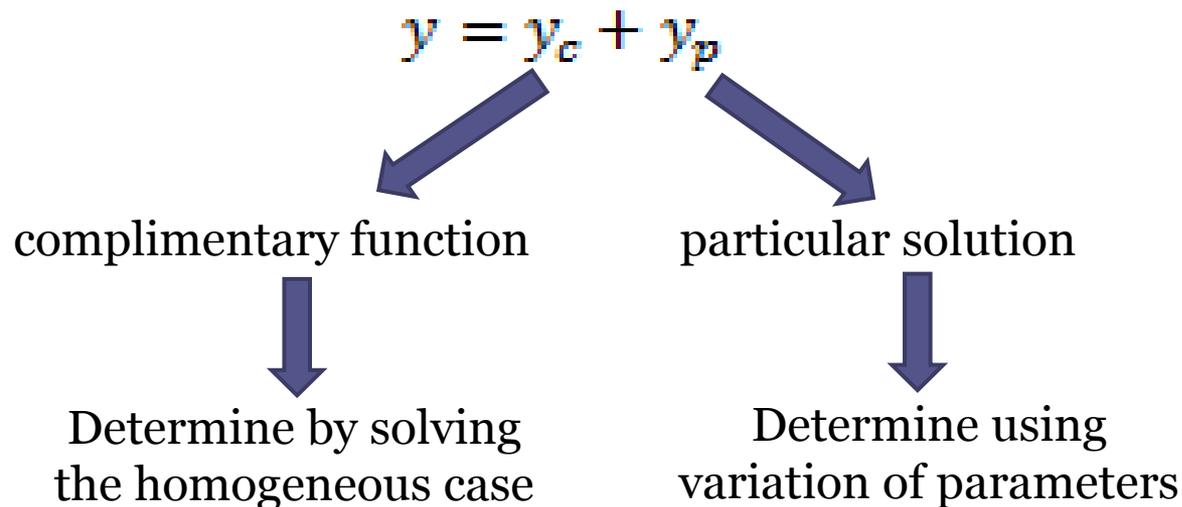
$$x^4 y'' + 2x^3 y' = c$$

Solving the Bioheat Equation

- The result is a second-order, linear, non-homogeneous differential equation

$$x^2 y'' + 2xy' = \frac{c}{x^2}$$

- The general solution to this form of DE is



The Complimentary Function

- The homogeneous case is: $x^2y'' + 2xy' = 0$
- This is a Cauchy-Euler equation, so the solution is of the form $y = x^m$
- Plug in y' and y'' to the homogeneous case and solve for m :

$$x^2m(m-1)x^{m-2} + 2xm x^{m-1} = 0$$



$$m = 0 \quad m = -1$$

- Since $y = x^m$, the complimentary function is given by:

$$y_c = c_1y_1 + c_2y_2$$

$$y_c = \alpha x^0 + \beta x^{-1}$$

$$y_c = \alpha + \frac{\beta}{x}$$

The Particular Solution

- To use variation of parameters, rewrite as:

$$y'' + \frac{2}{x}y' = \frac{c}{x^4}$$

- The solution is given by $y_p = u_1y_1 + u_2y_2$, where y_1 and y_2 are from the complimentary function and

$$u_1 = \int u_1' = \int \frac{W_1}{W} \quad u_2 = \int u_2' = \int \frac{W_2}{W} \quad W = \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$$

$$W_1 = \begin{vmatrix} 0 & y_2 \\ f(x) & y_2' \end{vmatrix} \quad W_2 = \begin{vmatrix} y_1 & 0 \\ y_1' & f(x) \end{vmatrix} \quad f(x) = \frac{c}{x^4}$$



$$y_p = \frac{c}{2x^2}$$

The General Solution

- The general solution is the sum of the complimentary function and the particular solution:

$$y = y_c + y_p = \alpha + \frac{\beta}{x} + \frac{c}{2x^2}$$

$$T(r) = \alpha + \frac{\beta}{r} + \frac{c}{2r^2}$$

$$T(r) = \alpha + \frac{\beta}{r} - \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r^2}$$

- Plugging this back into the bioheat equation verifies that it is a solution

The General Solution

- The solution is also valid in terms of units

Quantity	Unit
I	A
k	A/V*m
σ	V*A/m*K
α	K
β	K*m

- Need to determine α and β

Determining α

- Assume the tissue is unaffected by heating at an infinite distance from the electrode
 - As $r \rightarrow \infty$, $T(r) \rightarrow 310.15$ K (body temperature)

$$T(r) = \alpha + \frac{\beta}{r} - \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r^2} \longrightarrow \alpha = 310.15 \text{ K}$$

$$T(r) = 310.15 + \frac{\beta}{r} - \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r^2}$$

Determining β

- To determine a numeric value for β , we will need another boundary condition
 - Allow heat to exit system (otherwise temperature will rise indefinitely)
 - Assume heat cannot leave system at $r=r_0$
 - Use Fourier's Law: $\mathbf{q} = -k\nabla T$ (\mathbf{q} =local heat flux)
 - If no heat can flow at $r=r_0$, then:

$$0 = \nabla T$$

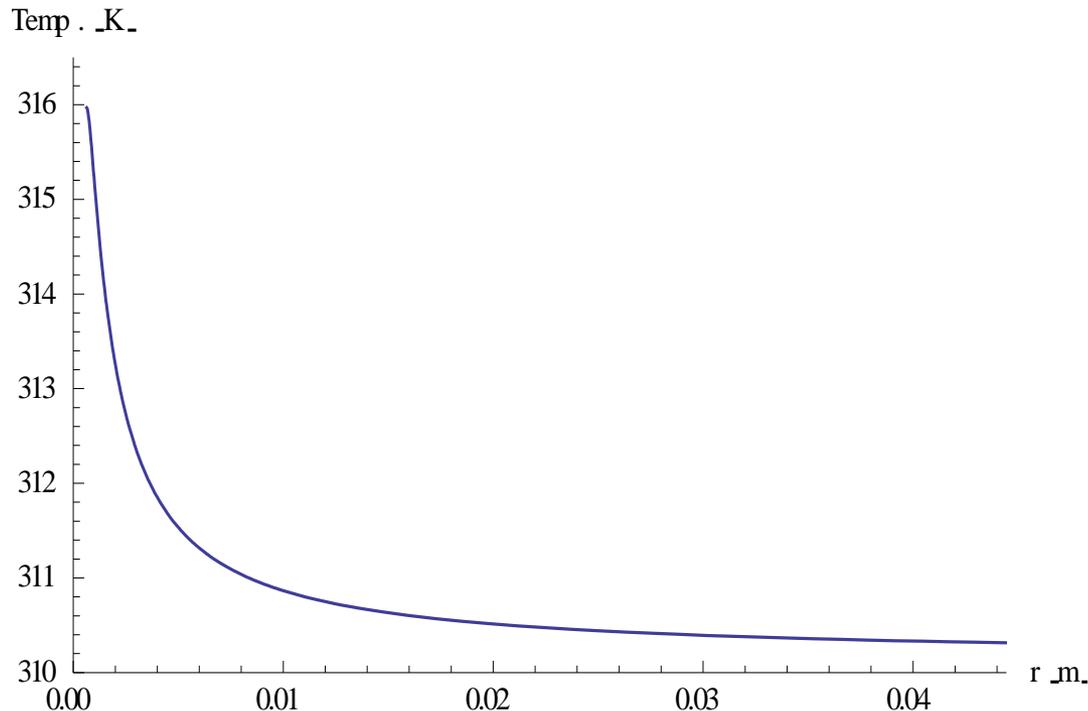
$$0 = -\frac{\beta}{r^2} - \frac{c}{r^3}$$

$$\beta = \left(\frac{I}{4\pi}\right)^2 \frac{1}{k\sigma r_0}$$

Determining β

- The solution for temperature is:

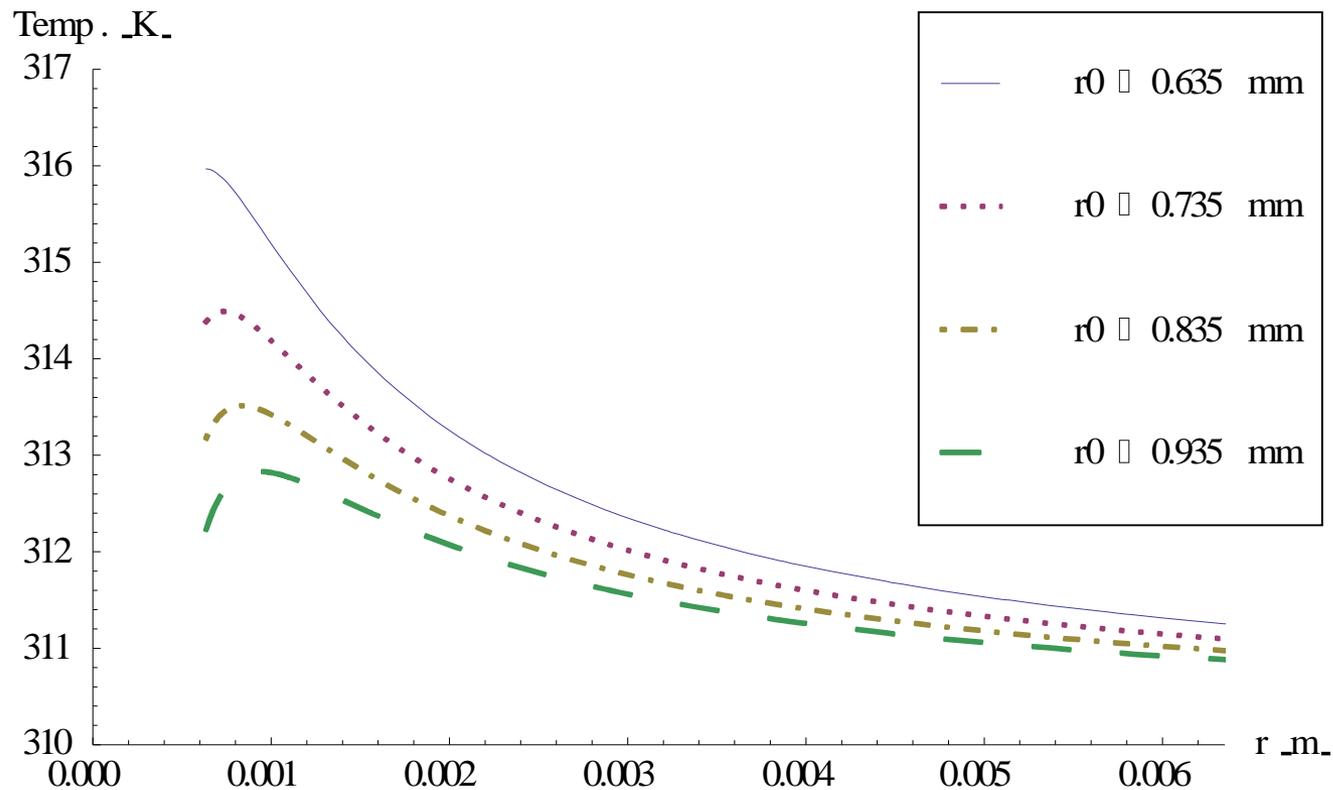
$$T(r) = 310.15 + \left(\frac{I}{4\pi}\right)^2 \frac{1}{k\sigma r_0} \frac{1}{r} - \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r^2}$$



A plot of temperature vs. radial distance from electrode ($I = 11.7$ mA, $\sigma = 0.327$ A/V*m, $k = 0.565$ W/m*K, $r_0 = 0.635$ mm).

Sensitivity to r_0

- Behavior of solution is highly dependent on r_0



Dependence of temperature on r_0

Future Work

- What is the *physical* meaning of the solution?
- The temperature distribution is

$$T(r) = 310.15 + \left(\frac{I}{4\pi}\right)^2 \frac{1}{k\sigma r_0} \frac{1}{r} - \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r^2}$$

or

$$T(r) = 310.15 + \frac{A}{r_0 r} - \frac{A}{2r^2} \quad (\text{where } A = \left(\frac{I}{4\pi}\right)^2 \frac{1}{k\sigma})$$

- What does it mean to have two similar terms competing?

Conclusions

- Recharging a subcutaneous medical device using electric fields can increase tissue temperature
- We show that the steady-state temperature distribution is given by $T(r) = 310.15 + \frac{\beta}{r} - \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r^2}$
- Future work: investigate how physical parameters influence temperature increase

Acknowledgments

- Wittenberg University
 - Dr. Daniel Fleisch
 - Dr. Elizabeth George
 - Dr. Adam Parker

- Duke University
 - Tom Jochum
 - Zachary Abzug
 - Dr. Patrick Wolf

References

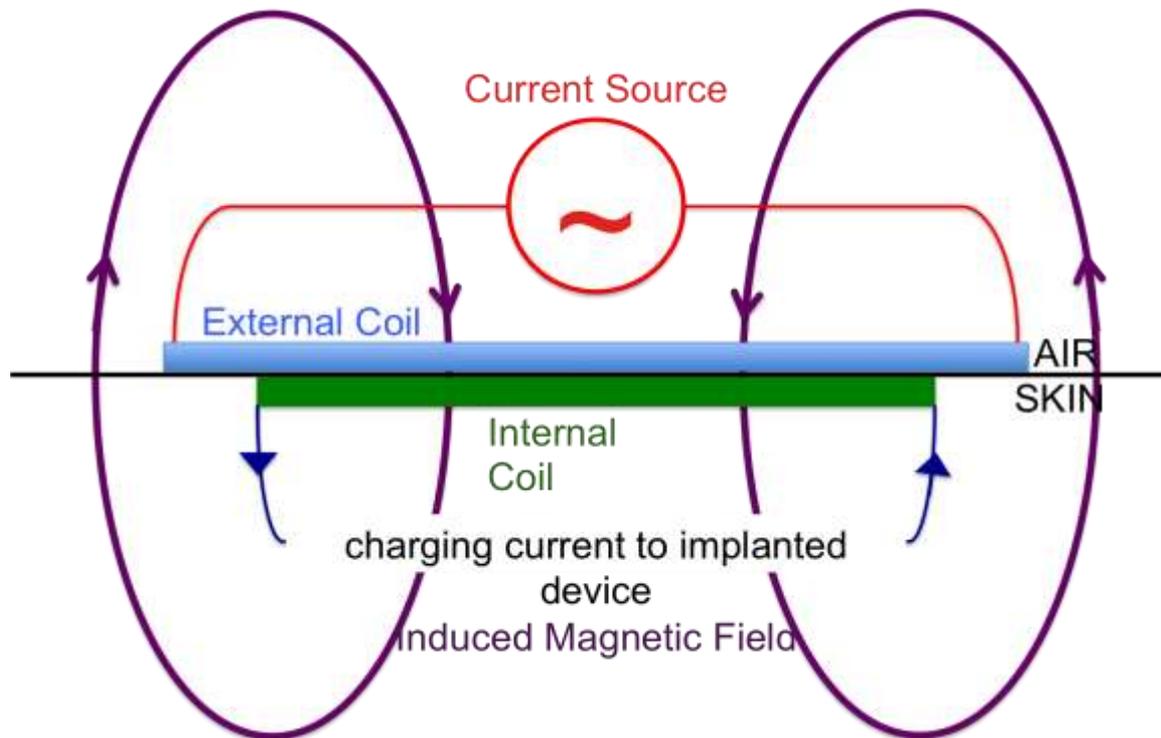
- [1] Elwassif, Maged M., Qingjun Kong, Maribel Vazquez, and Marom Bikson, "Bio-heat transfer model of deep brain stimulation-induced temperature changes," *J. Neural Eng.* **3**, 306-315, Nov. 2006
- [2] Griffiths, David J. "Conservation Laws." *Introduction to Electrodynamics*. Upper Saddle River, NJ: Prentice Hall, 1999. 345-46. Print.
- [3] Malmivuo, Jaakko and Robert Plonsey. *Bioelectromagnetism: Principles and Applications of Bioelectric and Biomagnetic Fields*. New York: Oxford University Press, 1995. Bioelectromagnetism. N.p., n.d. Web. 2 April 2012. <http://www.bem.fi/book/index.htm>.
- [4] Wissler, Eugene H., "Pennes' 1948 paper revisited," *J. Appl. Physiol.* **85**, 35-41, 1998

Questions?

Extra Slides

Transcutaneous Recharging

- Most implanted devices recharged via magnetic fields—not feasible for this device



Transcutaneous recharging using magnetic fields

Figure created by Zachary Abzug

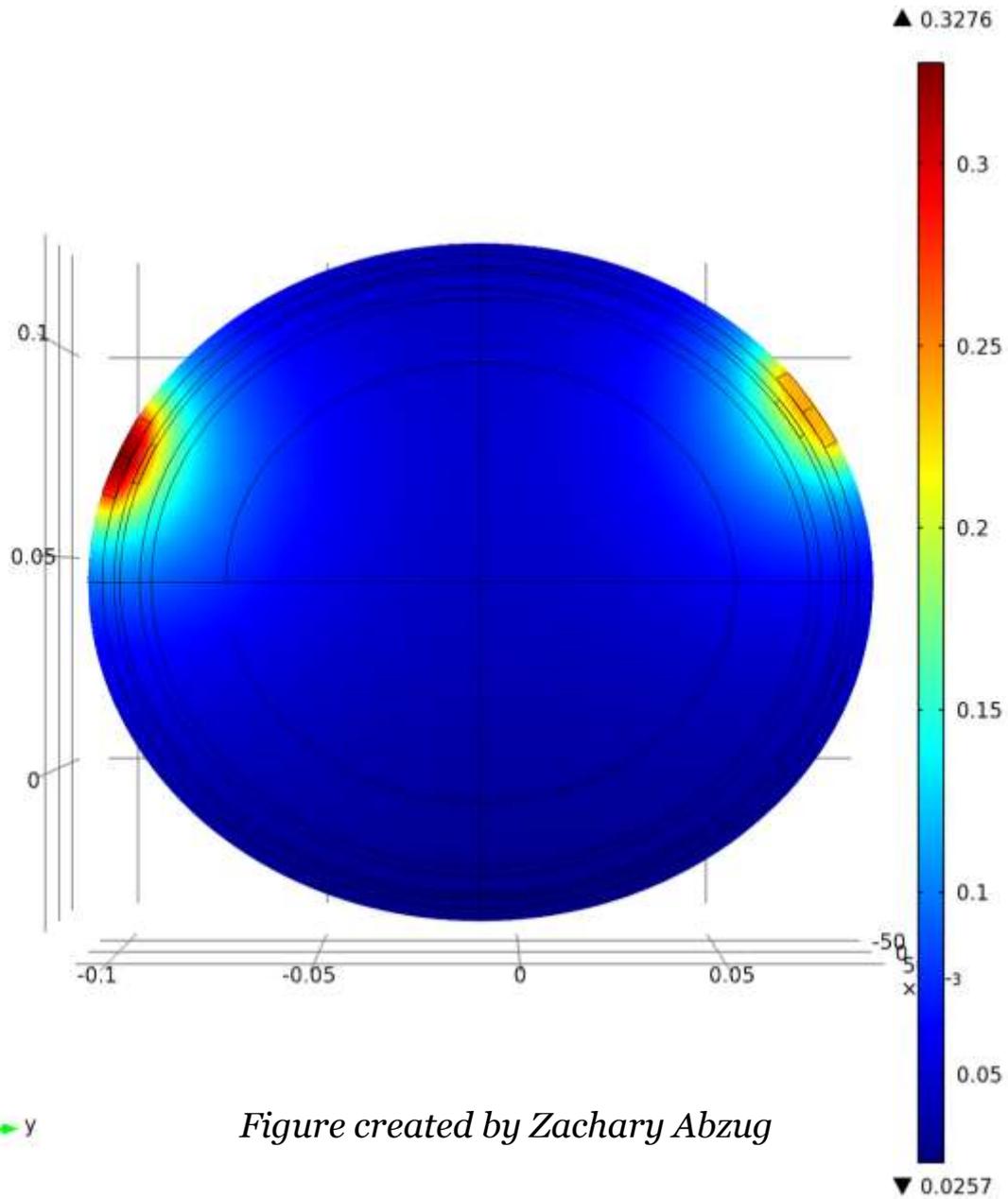


Figure created by Zachary Abzug

The Heat Equation

- The heat equation describes the three-dimensional variation of temperature in a region as a function of time

$$\rho C \frac{\partial T}{\partial t} = k \nabla^2 T$$

ρ = density

C = specific heat

k = thermal conductivity

- Not a complete model of heat transfer in biological situations due to perfusion (blood flow)

The Bioheat Equation

- The rate of heat transfer between blood and tissue is proportional to:
 - The volumetric perfusion rate
 - The difference between the arterial blood temperature and the local temperature
- Also add term (Q_{met}) to account for metabolic heat production
- The bioheat equation is:

$$\rho C \frac{\partial T}{\partial t} = k \nabla^2 T + \rho_b C_b \omega_b (T_b - T) + Q_{met}$$

ρ_b = density of blood

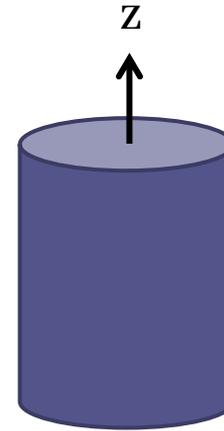
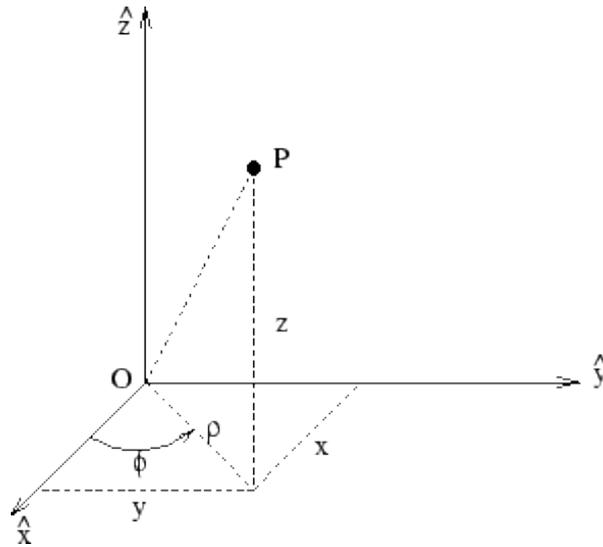
C_b = specific heat of blood

T_b = temperature of blood

ω_b = perfusion rate per unit volume of tissue

T = local tissue temperature

A Solution in Cylindrical Coordinates



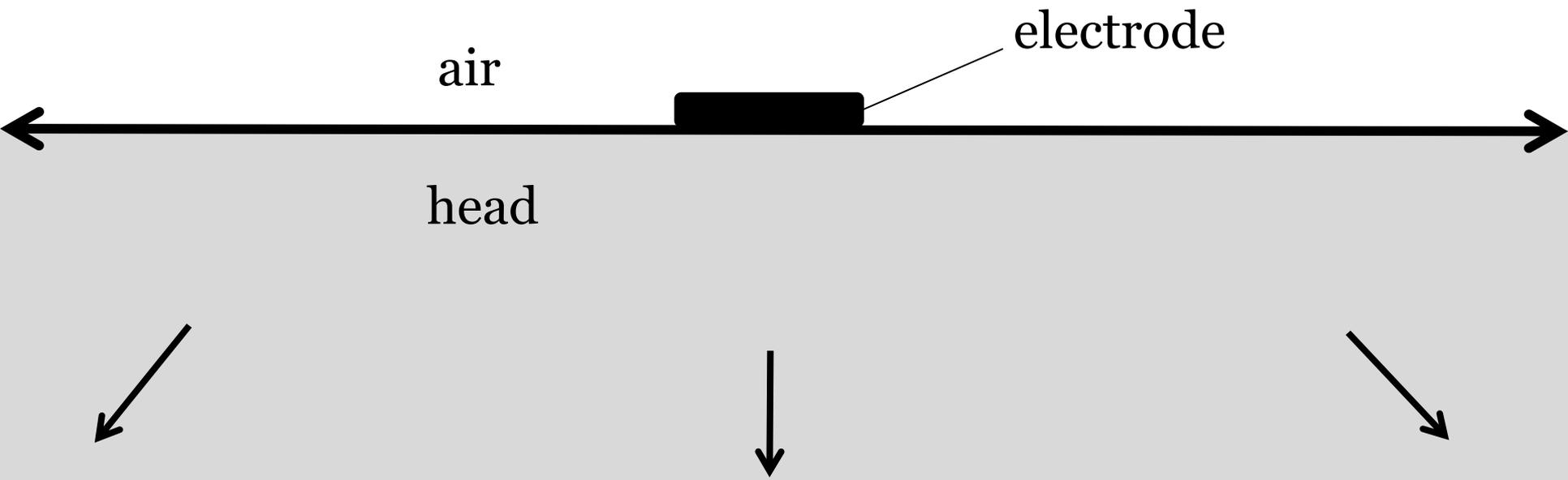
Cylindrical coordinate system
(image from uic.edu)

- Wrote Laplacian in cylindrical coordinates, ignoring ϕ and z dependence due to geometry

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) = -\frac{\sigma}{k} |\nabla V|^2$$

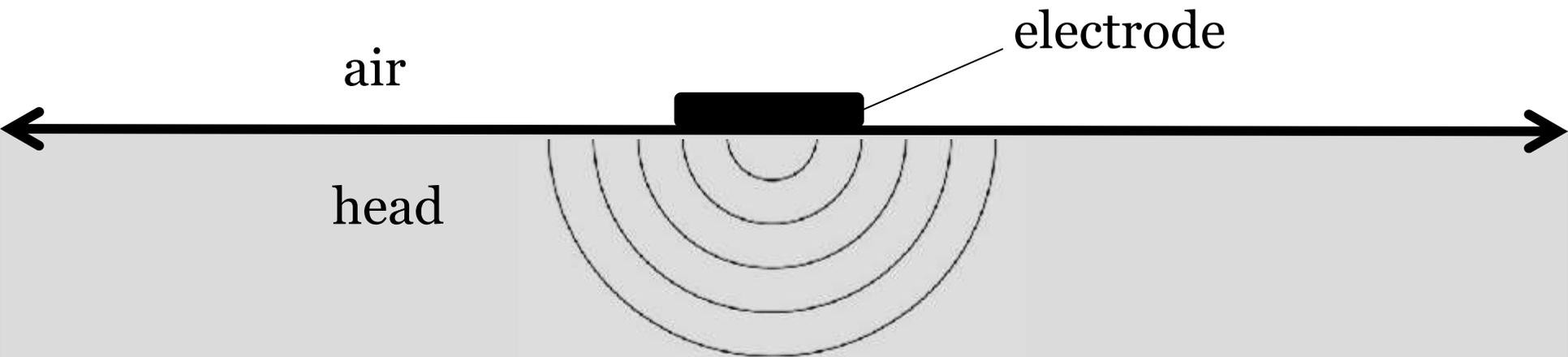
A Solution in Cylindrical Coordinates

- For our analysis, model head as infinitely wide and deep homogeneous resistive material
- Place one electrode on surface ($V=V_{\text{applied}}$) and one electrode at infinity ($V=0$)



A Solution in Cylindrical Coordinates

- It's acceptable to ignore φ dependence in our situation because of the axial symmetry
- Problem: we cannot ignore z dependence

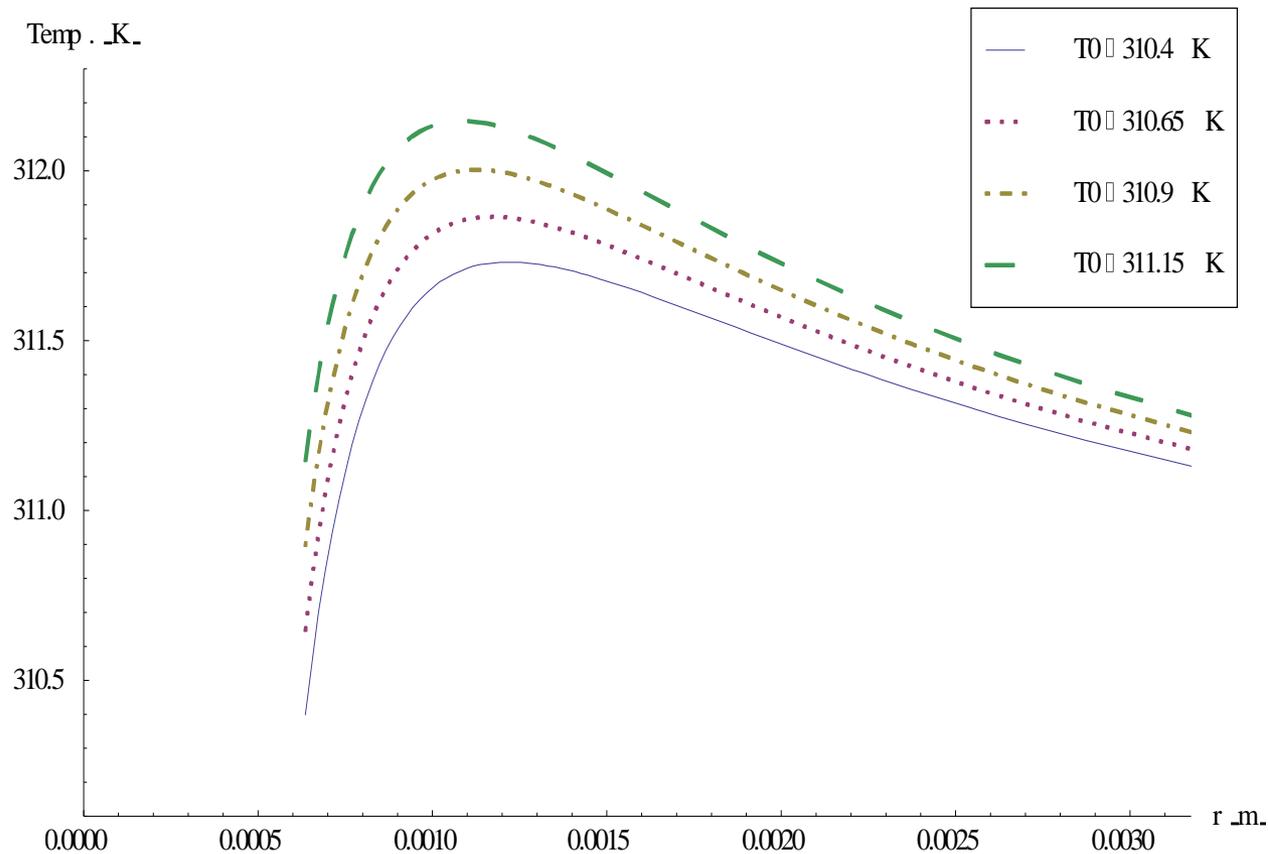


Determining β - Approach #1

- Want to know how the temperature behaves at the electrode's surface ($r=r_0$)
 - Requires knowing I , σ , k , r_0 , and T_0 (the temperature at the electrode's surface)
 - Choose representative values for I , σ , k , and r_0
 - Parameterize β based on selected values T_0
- Solve for β at r_0 : $\left(T_0 - 310.15 + \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r_0^2}\right) r_0 = \beta$
- The solution for temperature:

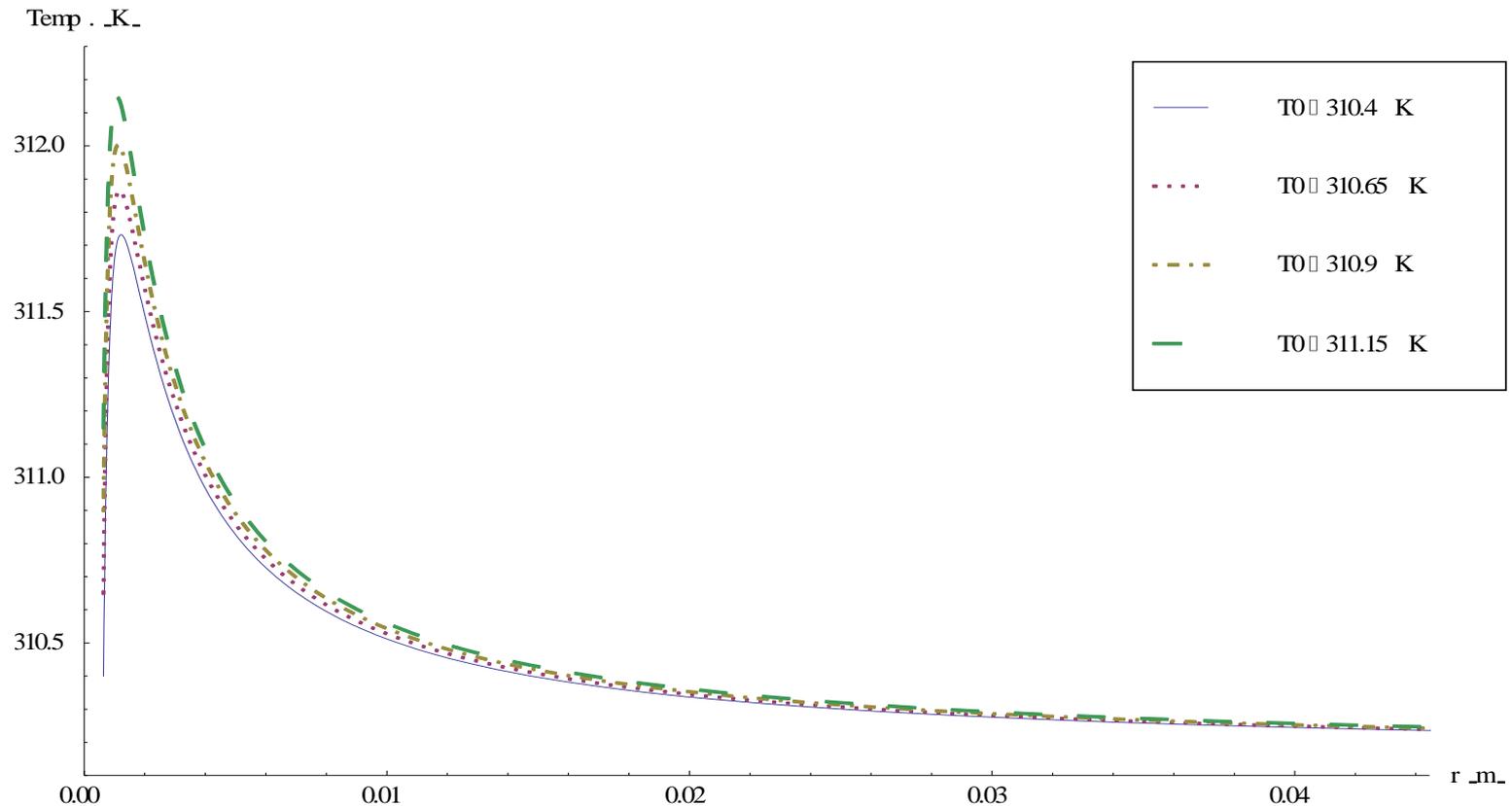
$$T(r) = 310.15 + \left(T_0 - 310.15 + \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r_0^2}\right) \frac{r_0}{r} - \left(\frac{I}{4\pi}\right)^2 \frac{1}{2k\sigma r^2}$$

Determining β - Approach #1



A plot of temperature vs. radial distance from electrode ($I = 11.7$ mA, $\sigma = 0.327$ A/V*m, $k = 0.565$ W/m*K, $r_o = 0.635$ mm).

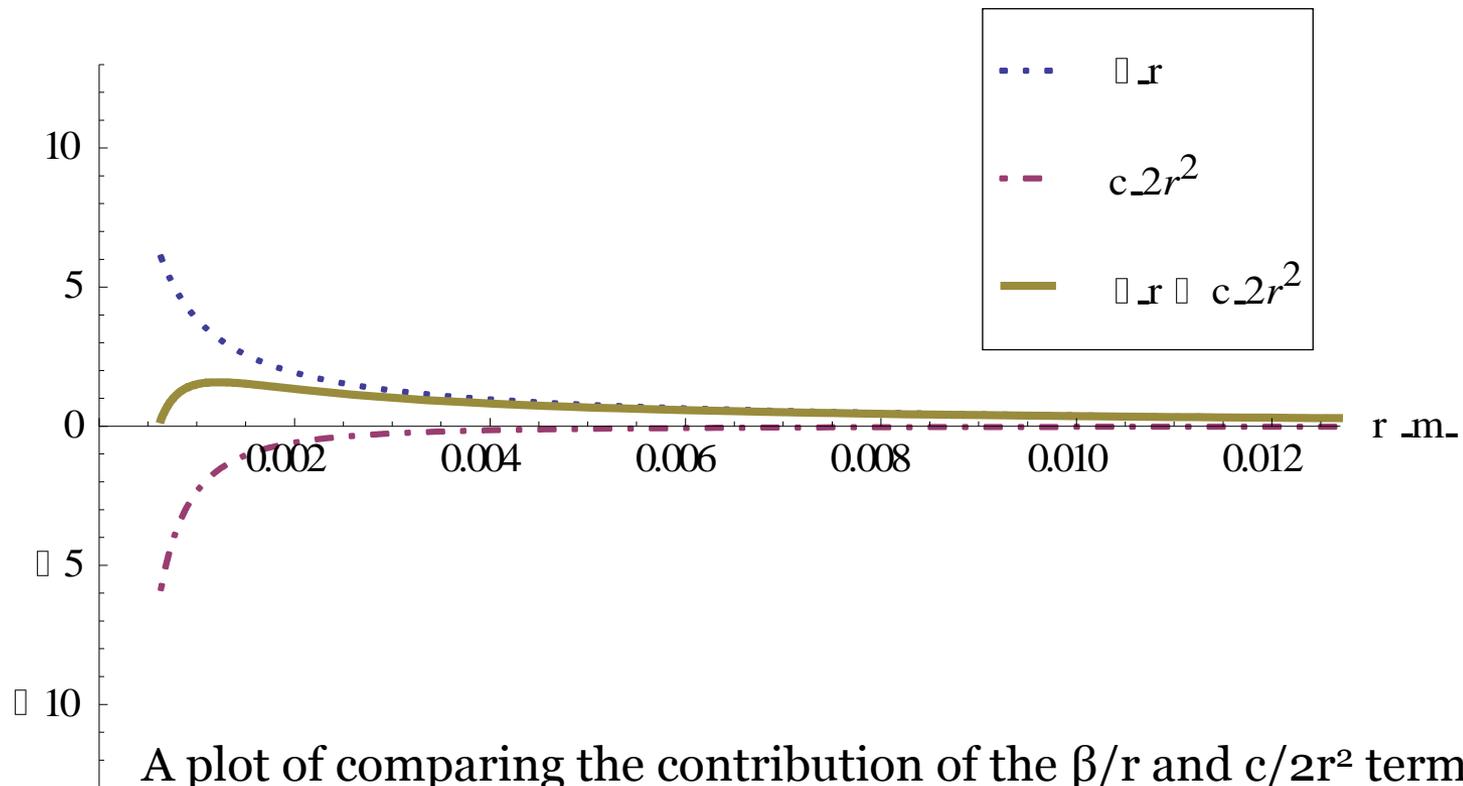
Determining β - Approach #1



A plot of temperature vs. radial distance from electrode ($I = 11.7 \text{ mA}$,
 $\sigma = 0.327 \text{ A/V}\cdot\text{m}$, $k = 0.565 \text{ W/m}\cdot\text{K}$, $r_o = 0.635 \text{ mm}$).

Temperature Peak

- If β is a smaller value, we see an initial peak in temperature
- Remember solution is of the form: $T(r) = \alpha + \frac{\beta}{r} + \frac{c}{2r^2}$

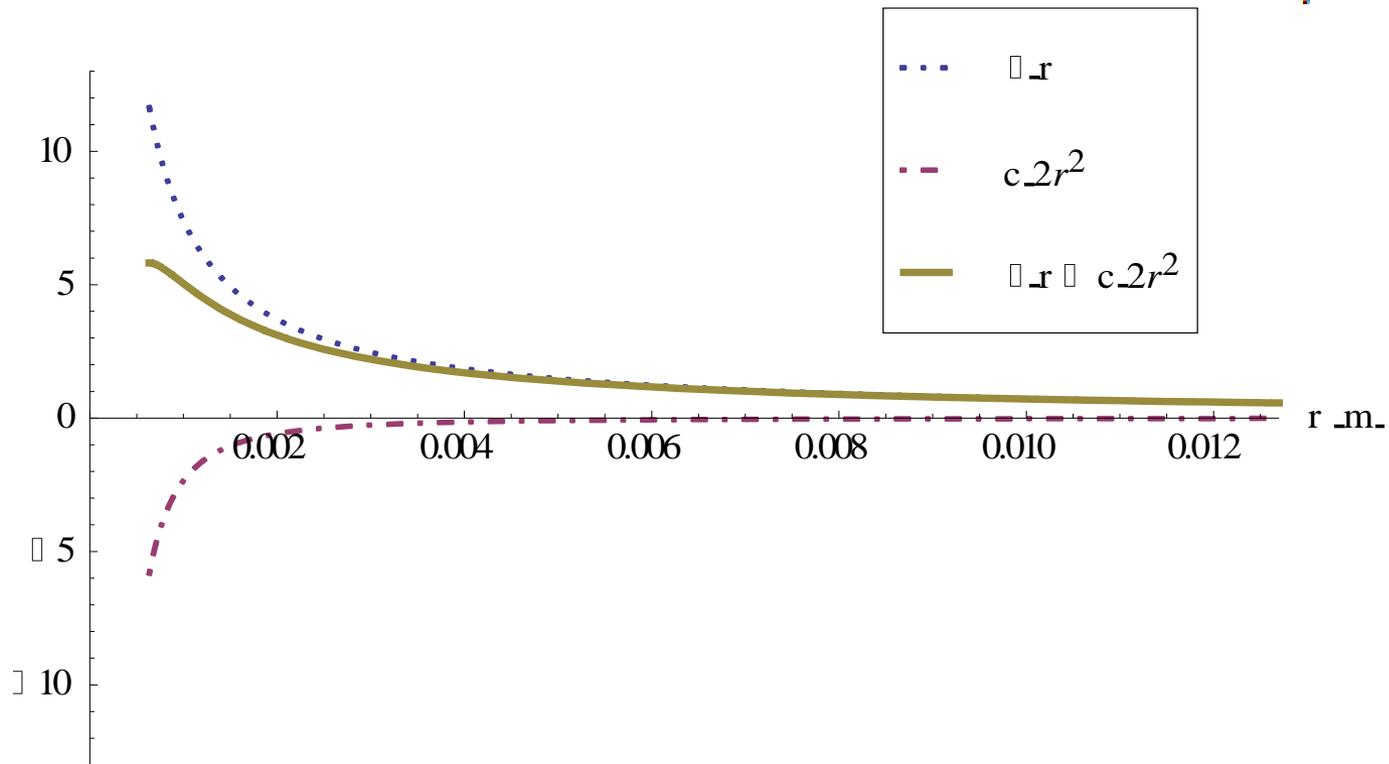


A plot of comparing the contribution of the β/r and $c/2r^2$ terms

Temperature Peak

- Would like to account for the peak in temperature

- Remember solution is of the form: $T(r) = \alpha + \frac{\beta}{r} + \frac{c}{2r^2}$



A plot of comparing the contribution of the β/r and $c/2r^2$ terms